

Stochastic Mean Field (SMF) description

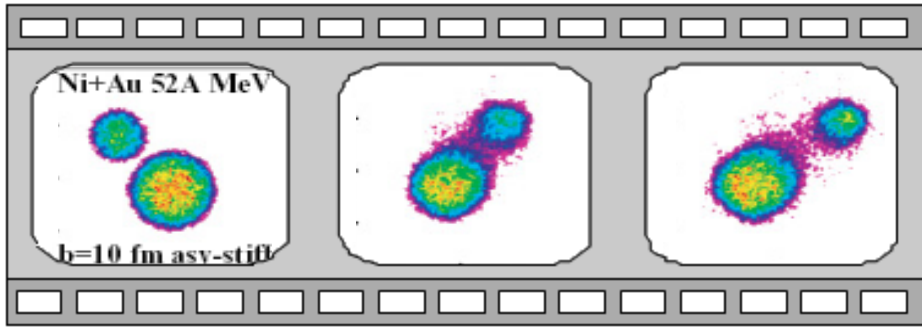
TRANSPORT 2017

March 27-30, 2017 FRIB-MSU, East Lansing, Michigan, USA

Maria Colonna

INFN - Laboratori Nazionali del Sud (Catania)

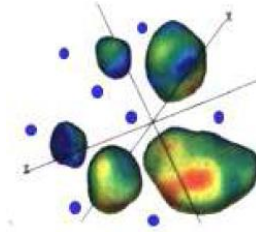
Dynamics of many-body system I



○ Mean-field (one-body) dynamics

○ Two-body correlations

○ Fluctuations



*one-body
density matrix*

two-body density matrix

$$\rho_2(12,1'2') = \underbrace{\rho_1(1,1')\rho_1(2,2')}_{\text{one-body}} + \delta\sigma(12,1'2')$$

$$H = H_0 + V_{1,2}$$

Mean-field

Residual interaction

$$i\hbar \frac{\partial}{\partial t} \rho_1(1,1',t) = \langle 1 | [H_0, \rho_1(t)] | 1' \rangle + K[\rho_1] + \delta K[\rho_1, \delta\sigma]$$

TDHF

$$K = F(\rho_1, |v|^2)$$

Average effect of the residual interaction

$$\delta K = F(v, \delta\sigma)$$

$$\langle \delta K \rangle = 0$$

$$\langle \delta K \delta K \rangle \rightarrow \text{Fluctuations}$$

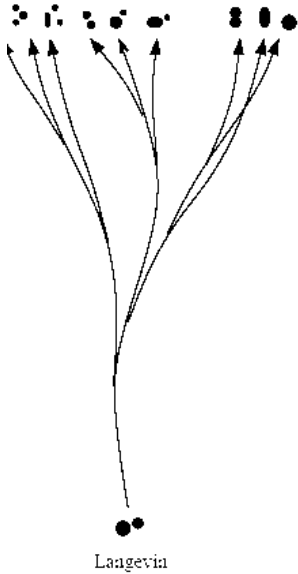
Dynamics of many-body systems II

Transition rate W
interpreted in terms of
NN cross section

Collision Integral

$$\mathbf{K} = g \sum_{234} W(12; 34) [\bar{f}_1 \bar{f}_2 f_3 f_4 - f_1 f_2 \bar{f}_3 \bar{f}_4]$$

$$\bar{f} = 1 - f$$



-- If statistical fluctuations larger than quantum ones

$$\langle \delta \mathbf{K}(p, t) \delta \mathbf{K}(p', t') \rangle = C \delta(t - t')$$

$$C(\mathbf{p}_a, \mathbf{p}_b, \mathbf{r}, t) = \delta_{ab} \sum_{234} W(a2; 34) F(a2; 34)$$

$$F(12; 34) \equiv f_1 f_2 \bar{f}_3 \bar{f}_4 + \bar{f}_1 \bar{f}_2 f_3 f_4.$$

Main ingredients:

- Residual interaction (2-body correlations and fluctuations)
- In-medium nucleon cross section

- Effective interaction
(self consistent mean-field)

Skyrme, Gogny forces

- Effective interactions

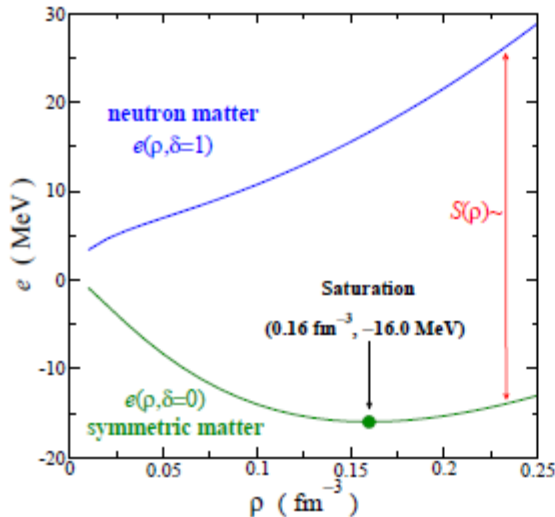
$$E = \langle \Psi | \hat{H} | \Psi \rangle$$

$$\approx \langle \Phi | \hat{H}_{eff} | \Phi \rangle = E[\hat{\rho}]$$

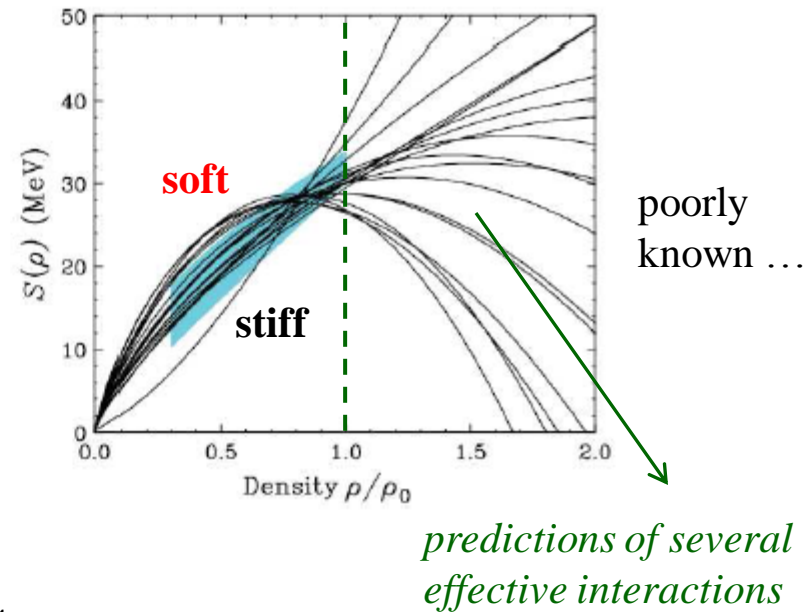
Energy Density Functional theories: The exact density functional is approximated with powers and gradients of one-body nucleon densities and currents.

The nuclear Equation of State ($T = 0$)

Energy per nucleon E/A (MeV)



Symmetry energy E_{sym} (MeV)



$$\frac{E}{A}(\rho, \beta) = \frac{E}{A}(\rho, \beta = 0) + E_{\text{sym}}(\rho)\beta^2 + O(\beta^4)$$

symm. matter

symm. energy

expansion around normal density

$$\beta = \text{asymmetry parameter} = (\rho_n - \rho_p)/\rho$$

➤ analogy with **Weizsacker mass formula** for nuclei (symmetry term) !

$$E_{\text{sym}}(\rho) = S_0 + L \frac{\rho - \rho_0}{3\rho_0} + \dots$$

(or J)

$$25 \leq J \leq 35 \text{ MeV} \quad 20 \leq L \leq 120 \text{ MeV}$$

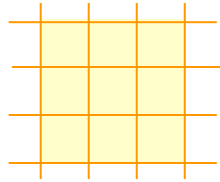
➤ 1. Semi-classical approximation to Nuclear Dynamics

Chomaz, Colonna, Randrup
 Phys. Rep. 389 (2004)
 Baran, Colonna, Greco, Di Toro
 Phys. Rep. 410, 335 (2005)

Transport equation for the one-body distribution function f

Semi-classical analog of the Wigner transform of the one-body density matrix

Density $f = f(\mathbf{r}, \mathbf{p}, t)$



Phase space (\mathbf{r}, \mathbf{p})

$$\frac{df(\mathbf{r}, \mathbf{p}, t)}{dt} = \frac{\partial f(\mathbf{r}, \mathbf{p}, t)}{\partial t} + \{f, H_0\} = 0$$

$$H_0 = T + U$$

Vlasov Equation,
 like Liouville equation:
 The phase-space density is
 constant in time

The mean-fields potential U is self-consistent: $U = U(\rho)$
 Nucleons move in the field created by all other nucleons

Semi-classical approximation \longrightarrow transport theories

$$\frac{df(\mathbf{r}, \mathbf{p}, t)}{dt} = \frac{\partial f(\mathbf{r}, \mathbf{p}, t)}{\partial t} + \{f, h\} = k[f] + \delta k$$

Vlasov

 $k[f] + \delta k$
Boltzmann-Langevin

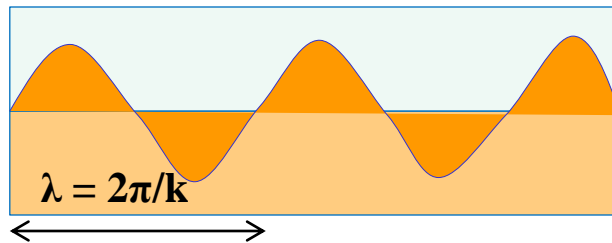
Correlations,
 Fluctuations

From BOB to SMF

- Fluctuations from *external stochastic* force (tuning of the most unstable modes)

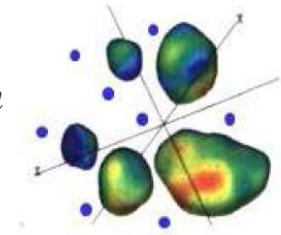
$$\dot{f} = \bar{I}[f] + \frac{\partial U_{\text{ext}}}{\partial r} \frac{\partial f}{\partial p}$$

Brownian One Body (**BOB**) dynamics



Chomaz, Colonna, Guarnera, Randrup
PRL73,3512(1994)

multifragmentation
event

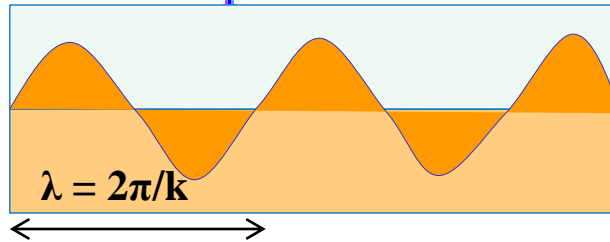


From BOB to SMF

- Fluctuations from *external* stochastic force (tuning of the most unstable modes)

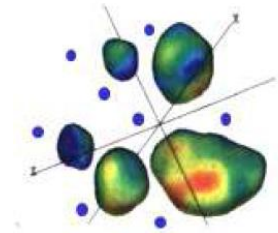
$$\dot{f} = \bar{I}[f] + \frac{\partial U_{\text{ext}}}{\partial r} \frac{\partial f}{\partial p}$$

Brownian One Body (**BOB**) dynamics

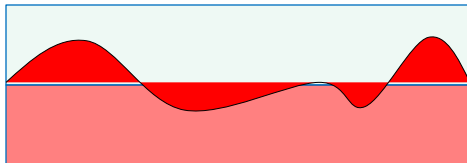


*Chomaz, Colonna, Guarnera, Randrup
PRL73,3512(1994)*

*multifragmentation
event*



- Stochastic Mean-Field (**SMF**) model :
Thermal fluctuations (at local equilibrium) are projected on the coordinate space by agitating the spacial density profile



M.Colonna et al., NPA642(1998)449

Details of the model

$$\triangleright f(\mathbf{r}, \mathbf{p}, t) = \frac{Ch^3}{4} \sum_i g_r(\mathbf{r} - \mathbf{r}_i) g_p(\mathbf{p} - \mathbf{p}_i),$$

triangular functions in \mathbf{r} space (for g_r) and δ functions in momentum space (for g_p)

$$\text{triangular function} \quad g(x^j - x_i^j) = 2l - |x^j - x_i^j|,$$

$$l = 1 \text{ fm}$$

Total number of test particles: $N_{\text{tot}} = N_{\text{test}} * A$ *lattice size*

\triangleright System total energy (lattice Hamiltonian):

$$E_{\text{tot}} = \sum_i p_i^2 / (2m) + N_{\text{test}} \left[\int d\mathbf{r} \rho(\mathbf{r}) E_{\text{pot}}(\rho_n, \rho_p) \right.$$

$$\left. + \int d\mathbf{r} \rho_p(\mathbf{r}) E_{\text{pot}}^{\text{Coul}}(\rho_p) / 2, \right]$$

➤ Potential energy

$$E_{pot}(\rho) = \frac{A}{2}\tilde{\rho} + \frac{B}{\sigma + 1}\tilde{\rho}^\sigma + \frac{C_{surf}}{2\rho}(\nabla\rho)^2 + \frac{1}{2}C_{sym}(\rho)\tilde{\rho}\beta^2,$$

where $\tilde{\rho} = \rho/\rho_0$ (ρ_0 denotes the saturation density),
 $A = -356 \text{ MeV}$, $B = 303 \text{ MeV}$, $\sigma = 7/6$. β = asymmetry parameter

➤ $C_{surf} = -6/\rho_0^{5/3} \text{ MeV fm}^5$.

Negative surface term to correct surface effects induced by the the use of finite width t.p. packets

➤ Symmetry energy parametrizations :

asy-EoS	E_{sym}/A	L(MeV)
asysoft	30.	14.
asystiff	28.	73.
asysupstiff	28.	97.

New Skyrme interactions
 (SAMi-J family) recently introduced

Hua Zheng et al.

PHYSICAL REVIEW C 94, 014313 (2016)

- the Coulomb potential is determined solving the Laplace equation:

$$\nabla^2 E_{pot}^{Coul} = -4\pi e^2 \rho_p = -18.1 \rho_p$$

→ *Initialization and dynamical evolution*

- Ground state initialization with Thomas-Fermi
- Test particle positions and momenta are propagated according to the Hamilton equations (non relativistic)

Details of the model: Collision Integral

➤ Mean free path method :

each test particle has just one collision partner

$$\tau_{col} = \frac{\lambda}{v_{kl}} = \frac{1}{\rho \sigma_{NN} v_{kl}}, \quad P_{col}(\Delta t) \approx \Delta t / \tau_{col}$$

$\Delta t = \text{time step}$

➤ blocking factors, defined as $P_{Pauli} = (1 - f_l)(1 - f_k)$,

we now take a Θ function in \mathbf{r} space and a gaussian function, with $\sigma = 29 \text{ MeV}/c$, in momentum space. The Θ function is defined as: $\Theta(\mathbf{r} - \mathbf{r}_i) = \Theta(R - |\mathbf{r} - \mathbf{r}_i|)$, with $R = 2.53 \text{ fm}$. The new definition makes the occupation number smoother (though less local), reducing fluctuations which may induce spurious collisions.

➤ Free n-p and p-p cross sections, with a maximum cutoff of 50 mb

Fluctuation tuning

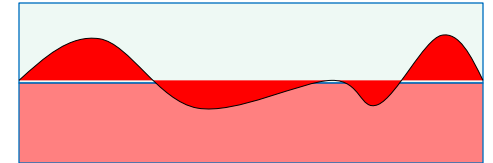
When local equilibrium is achieved:

➤ $\sigma_f^2 = f(1-f)$ → Fluctuation variance for a fermionic system at equilibrium

$$\sigma_\rho^2(\mathbf{r}, t) = \frac{1}{V} \int \frac{d\mathbf{p}}{h^3/4} \sigma_f^2(\mathbf{r}, \mathbf{p}, t).$$

$$\sigma_\rho^2 = \frac{\rho}{V} \frac{3T}{2\varepsilon_F} \left[1 - \frac{\pi^2}{12} \left(\frac{T}{\varepsilon_F} \right)^2 + \dots \right].$$

➤ Stochastic Mean-Field (**SMF**) model :
Fluctuations are projected on the coordinate space by agitating the spacial density profile



- [1] V. Baran, M. Colonna, V. Greco and M. Di Toro, Phys. Rep. **410**, 335 (2005)
- [2] M. Colonna et al., Nucl. Phys. A **642**, 449 (1998)

Ph. Chomaz, M. Colonna, J. Randrup, Phys. Rep. **389**, 263 (2004)

Some applications

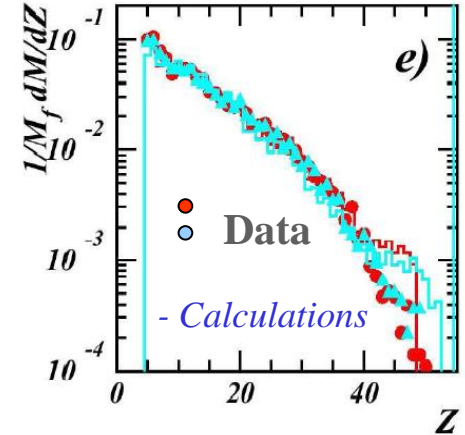
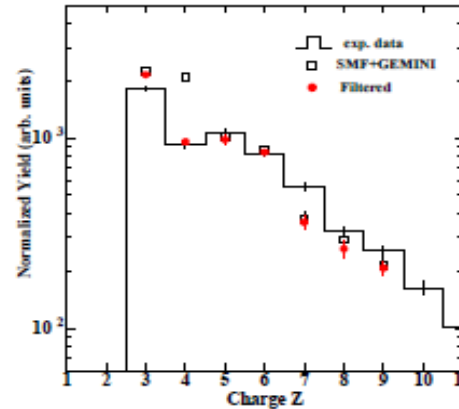
Fragmentation studies
in central and semi-peripheral collisions

Isospin effects at Fermi energies

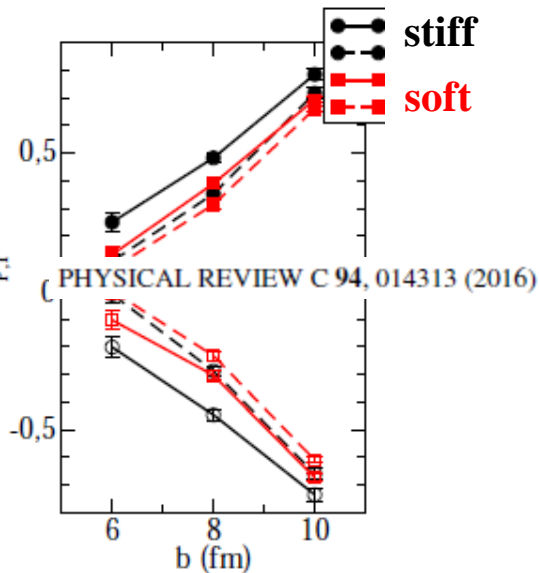
Small amplitude dynamics (collective modes)
and low-energy
reaction dynamics

Charge distribution

E. De Filippo et al., PRC(2012)



J. Frankland et al.,
NPA 2001



J. Rizzo et al., NPA(2008)

Hua Zheng et al.

PHYSICAL REVIEW C 94, 014313 (2016)

